

Capítulo 12

1. Exercício 16.1: temos a condição inicial $C\Delta\varphi = Q_0$, que é a carga inicial no capacitor, após ter sido carregado. Também temos que $\mathcal{J} = \rho\mathcal{E}$. Usando a equação da continuidade, dada por

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathcal{J} = 0$$

ficamos com

$$\frac{\partial \rho}{\partial t} + \rho \nabla \cdot \mathcal{E} = 0$$

e, usando a eq. de Maxwell $\nabla \cdot \mathcal{E} = \frac{\rho}{\epsilon_0}$, ficamos com (tendo integrado no volume)

$$\frac{\partial q(t)}{\partial t} = -\frac{\rho}{\epsilon_0} q(t)$$

o que dá, após integração e substituição da condição inicial, o resultado

$$q(t) = C\Delta\varphi e^{-t/\tau},$$

como desejado. O item (b) pode ser obtido imediatamente a partir da derivação do resultado anterior e é dado por

$$i(t) = -\frac{\partial}{\partial t} C\Delta\varphi e^{-t/\tau}.$$

Finalmente, o item (c) é imediatamente obtido como $\mathcal{B} = 0$ pela simetria do problema.

2. Exercício 16.3: O item (a) pode ser obtido se usarmos a equação de Maxwell

$$\nabla \times \mathcal{H} = \frac{\partial \mathcal{D}}{\partial t} + \mathcal{J}$$

e a integramos sobre a superfície do dielétrico. A integral de superfície de $\nabla \times \mathcal{H}$ torna-se uma integral de linha sobre \mathcal{H} e a integral de superfície sobre \mathcal{D} é simples de ser obtida, já que \mathcal{D} é constante. Assim, ficamos com

$$\oint_C \mathcal{H} \cdot d\mathcal{L} = \frac{\partial \mathcal{D}}{\partial t} \cdot \iint_S n dA,$$

onde já usamos o fato de que, no dielétrico, $\mathcal{J} = 0$. Assim, ficamos com (note que \mathcal{H} é constante sobre C , por simetria)

$$2\pi a H_E = \pi a^2 \frac{\partial D_z}{\partial t},$$

onde já expressamos o resultado em termos das componentes distintas de zero dos campos. Assim,

$$H_E = \frac{1}{2} a \frac{\partial D_z}{\partial t}.$$

Reitz Electromagnetic Solutions

Michal Rosen-Zvi



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